

Global Regularity for Very Weak Solutions to Boundary Value Problem of Homogeneous A -Harmonic Equation

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Abstract—The very weak solution to elliptic boundary value problems is considered. A global regularity result is derived for very weak solutions under some controllable and coercivity conditions, by using the Hodge decomposition theorem and the methods in Sobolev spaces.

Keywords—Hodge decomposition theorem; A -harmonic equation; global regularity

I. INTRODUCTION

Let Ω be a bounded regular open set of \mathbf{R}^n ($n \geq 2$). We consider the boundary value problem for the second order degenerate elliptic equation

$$\begin{cases} \operatorname{div} A(x, \nabla u) = 0, & \text{in } \Omega \\ u = u_0, & \text{on } \partial\Omega \end{cases} \quad (1.1)$$

where $A(x, \xi): \Omega \times \mathbf{R}^n \rightarrow \mathbf{R}^n$ is a Carathéodory function satisfying the coercivity and growth conditions: for almost all $x \in \Omega$, all $\xi \in \mathbf{R}^n$,

$$(H1) \quad |A(x, \xi)| \leq \beta |\xi|^{p-1},$$

$$(H2) \quad \langle A(x, \xi), \xi \rangle \geq \alpha |\xi|^p,$$

where $1 < p < \infty$, $0 < \alpha \leq \beta < \infty$, $u_0 \in W^{1,r}(\Omega)$ is a boundary value function.

Definition 1.1 A function $u \in u_0 + W_0^{1,r}(\Omega)$, $\max\{p, -\} \leq r < p$ is called a very weak solution to (1.1), if

$$\int_{\Omega} \langle A(x, \nabla u), \nabla \phi \rangle dx = 0$$

holds true for any $\phi \in W_0^{1, \frac{r}{r-p+1}}(\Omega)$ with compact support in Ω .

A crucial fact is that r can be smaller than the natural exponent p . For variational extremals the global higher integrability of the derivative ∇u has been studied by Granlund S^[1] in the case $p = n$. For this it

seems necessary to impose a regularity condition for $\partial\Omega$.

We say that $\partial\Omega$ is r -Poincaré thick, if there is $0 < C < \infty$ such that for all open cube $Q_R \subset \mathbf{R}^n$ with side length $R > 0$, it holds

$$\left(\int_{Q_{2R}} |u|^r dx \right)^{\frac{1}{r}} \leq C \left(\int_{Q_{2R}} |\nabla u|^{\frac{m}{r+n}} dx \right)^{\frac{r+n}{m}} \quad (1.2)$$

whenever $u \in W^{1,r}(Q_{2R})$, $u = 0$ a.e. on $(\mathbf{R}^n \setminus \Omega) \cap Q_{2R}$, and $Q_{\frac{3R}{2}} \cap \Omega^c \neq \emptyset$. Here, and in the following, $Q(\lambda R)$, $\lambda > 0$, means a cube parallel to $Q(R)$ with the same center as $Q(R)$ and with side length λR . See [2].

The following is the main conclusion of this paper.

Theorem 1.2 Suppose that a bounded regular domain Ω has a r -Poincaré thick boundary and that $r \geq \frac{n}{n-1}$, operator A satisfy conditions (H1)-(H2). If $u_0 \in W^{1,r}(\Omega)$ is the boundary value function, $u \in W^{1,r}(\Omega)$ is the very weak solution of Dirichlet problem (1.1), then there exists $R_0 > 0$ and r_1, r_2 , satisfying

$$r_1 = r_1(n, p, K, R_0, \alpha, \beta, \Omega) < p < r_2 = r_2(n, p, K, R_0, \alpha, \beta, \Omega),$$

such that $\forall r \in [r_1, p)$, $u \in W^{1,r_2}(\Omega)$, then u is the weak solution in the classical meaning.

II. PRELIMINARY LEMMAS

Let Ω be a bounded regular domain, $x_0 \in \Omega$, $0 < R < \operatorname{dist}(x_0, \partial\Omega)$, $Q_R(x_0) \subset \Omega$, here $Q_R(x_0)$ is a cube with side length of R and a center of x_0 .

Lemma 2.1 [3] Let $1 < p < n$, $0 < q \leq \frac{np}{n-p}$ if $u \in W^{1,p}(B_R(x_0))$, then

$$\|u - u_R\|_{L^q(B_R(x_0))} \leq CR^{\frac{n(\frac{1}{q} - \frac{1}{p}) + 1}{p}} \|\nabla u\|_{L^p(B_R(x_0))}, \quad (2.1)$$

here $u_R = \frac{1}{|B_R(x_0)|} \int_{B_R(x_0)} u dx$, C is a positive constant only depending on p, q, n .

Specially, if $u \in W_0^{1,p}(B_R(x_0))$, then

$$\|u\|_{L^q(B_R(x_0))} \leq CR^{\frac{n(\frac{1}{q} - \frac{1}{p}) + 1}{p}} \|\nabla u\|_{L^p(B_R(x_0))}. \quad (2.2)$$

This lemma gives the dependence of embedding theorem on region size. For the case $q = \frac{np}{n-p}$ see [4].

This lemma is a direct corollary of theorem 7.10 and the Hölder inequality in Gilbarg-Trudinger [4].

This lemma also applies to cubes.

Lemma 2.2 [3] (Hodge decomposition) Let $\Omega \subset \mathbf{R}^n$ be a regular domain and $u \in W_0^{1,r}(\Omega, \mathbf{R}^m)$, and let $0 < \varepsilon < r-1, r = p - \varepsilon \geq \max\{1, p-1\}$. Then there exist $\phi(x) \in W_0^{1, \frac{r}{1-\varepsilon}}(\Omega, \mathbf{R}^m)$ and a (divergence free) matrix-field $H(x) \in L^{\frac{r}{1-\varepsilon}}(\Omega, \mathbf{R}^{n \times m})$, such that

$$|\nabla u|^{-\varepsilon} \nabla u = \nabla \phi + H. \quad (2.3)$$

Moreover

$$\|H\|_{\frac{r}{1-\varepsilon}} \leq C\varepsilon \|\nabla u\|_r^{-\varepsilon} \quad (2.4)$$

where C is a constant that only depends on n, r and Ω .

Remark 2.3 It can be seen from (2.3) and (2.4), estimates of $\nabla \phi$ are similar to those of (2.4).

Lemma 2.4 [5] Suppose X and Y are vectors of an inner product space, $0 \leq \varepsilon < 1$. Then

$$\| |X|^{-\varepsilon} X - |Y|^{-\varepsilon} Y \| \leq \frac{2^\varepsilon (1+\varepsilon)}{1-\varepsilon} \|X - Y\|^{1-\varepsilon}.$$

Lemma 2.5 [6] (Reverse Hölder inequality) Let Q be an n -cube. Suppose

$$\int_{Q_R(x_0)} g^q dx \leq \theta \int_{Q_{2R}(x_0)} g^q dx + c \left(\int_{Q_{2R}(x_0)} g dx \right)^q + \int_{Q_{2R}(x_0)} f^q dx$$

for each $x_0 \in Q$ and each $R < \frac{1}{2} \text{dist}(x_0, \partial Q) \wedge R_0$, where

R_0, b, θ are constants with $b > 1, R_0 > 0, 0 \leq \theta < 1$. Then $g \in L_{loc}^p(Q)$ for $p \in [q, q + \varepsilon)$ and

$$\left(\int_{Q_R(x_0)} g^p dx \right)^{\frac{1}{p}} \leq c \left\{ \left(\int_{Q_{2R}(x_0)} g^q dx \right)^{\frac{1}{q}} + \left(\int_{Q_{2R}(x_0)} f^p dx \right)^{\frac{1}{p}} \right\}$$

for $Q_{2R} \subset Q, R < R_0$, where c and ε are positive constants only depending on b, θ, q, n .

III. PROOF OF THEOREM 1.2

Proof. Let $x_0 \in \Omega, Q_\rho = Q_\rho(x_0)$ is a cube with side length of ρ and a center of x_0 . ε is a sufficiently small positive number, $r = p - \varepsilon$. Since Ω is bounded, we can choose a cube $Q_0 = Q_{2R_0}$ such that $\Omega \subset Q_{R_0}$. Next let $Q_{2R} \subset Q_0$. There are two possibilities: (1) $Q_{\frac{3}{2}R} \subset \Omega$; (2) $Q_{\frac{3}{2}R} \cap \Omega^c \neq \emptyset$.

In the case (1), for $Q_{\frac{3}{2}R} \subset \Omega$, fix a cutoff function $\eta \in C_0^\infty(Q_{\frac{3}{2}R})$ such that $0 \leq \eta \leq 1, |\nabla \eta| \leq \frac{C}{R}$, and $\eta \equiv 1$ on $x \in Q_R$. Let $u \in W^{1,r}(\Omega)$ be a very weak solution of problem(1.1). Consider the following Hodge decomposition

$$|\nabla(\eta u)|^{-\varepsilon} \nabla(\eta u) = \nabla \phi + H, \quad (3.1)$$

here $\phi \in W_0^{1, \frac{r}{1-\varepsilon}}(Q_{\frac{3}{2}R}), H \in L^{\frac{r}{1-\varepsilon}}(Q_{\frac{3}{2}R})$ is a (divergence free) matrix-field, satisfying

$$\|\nabla \phi\|_{\frac{r}{1-\varepsilon}} \leq C \|\nabla(\eta u)\|_r^{-\varepsilon}, \quad (3.2)$$

$$PHP_{\frac{r}{1-\varepsilon}} \leq C\varepsilon P \nabla(\eta u) P_r^{-\varepsilon}. \quad (3.3)$$

Let

$$E(\eta, u) = |\nabla(\eta u)|^{-\varepsilon} \nabla(\eta u) - |\eta \nabla u|^{-\varepsilon} \eta \nabla u, \quad (3.4)$$

by Lemma 2.4 we have

$$|E(\eta, u)| \leq 2^\varepsilon \frac{1+\varepsilon}{1-\varepsilon} |\eta \nabla u|^{1-\varepsilon}. \quad (3.5)$$

A useful technique in the following calculation is to use ϕ in Hodge decomposition (3.1) as the test function in Definition 1.1. Then

$$\begin{aligned} & \int_{Q_{\frac{3}{2}R}} \langle A(x, \nabla u), E(\eta, u) \rangle dx \\ & + \int_{Q_{\frac{3}{2}R}} \langle A(x, \nabla u), |\eta \nabla u|^{-\varepsilon} \eta \nabla u \rangle dx \\ & = \int_{Q_{\frac{3}{2}R}} \langle A(x, \nabla u), H \rangle dx, \end{aligned} \quad (3.6)$$

that is

$$\begin{aligned} & \int_{Q_{\frac{3}{2}R}} \langle A(x, \nabla u), \eta \nabla |\eta \nabla u|^{-\varepsilon} \eta \nabla u \rangle dx \\ & = \int_{Q_{\frac{3}{2}R}} \langle A(x, \nabla u), H \rangle dx - \int_{Q_{\frac{3}{2}R}} \langle A(x, \nabla u), E(\eta, u) \rangle dx \\ & = I_1 + I_2. \end{aligned} \quad (3.7)$$

Let's first estimate the left side of formula (3.7). By the hypothesis condition (H2) and the definition of η ,

$$\begin{aligned} & \int_{Q_{\frac{3}{2}R}} \langle A(x, \nabla u), |\eta \nabla u|^{-\varepsilon} \eta \nabla u \rangle dx \\ &= \int_{Q_{\frac{3}{2}R}} \eta^{1-\varepsilon} |\nabla u|^{-\varepsilon} \langle A(x, \nabla u), \nabla u \rangle dx \\ &\geq \alpha \int_{Q_{\frac{3}{2}R}} \eta^{1-\varepsilon} |\nabla u|^{-\varepsilon} |\nabla u|^p dx \\ &\geq \alpha \int_{Q_R} |\nabla u|^r dx. \end{aligned} \tag{3.8}$$

The estimate of I_1 is given below. By the hypothesis (H1), the Hölder inequality and (3.3), we can get the result

$$\begin{aligned} |I_1| &= \left| \int_{Q_{\frac{3}{2}R}} \langle A(x, \nabla u), H \rangle dx \right| \leq \int_{Q_{\frac{3}{2}R}} |A(x, \nabla u)| |H| dx \\ &\leq \beta \int_{Q_{\frac{3}{2}R}} |\nabla u|^{p-1} |H| dx \\ &\leq \beta \|\nabla u\|_r^{p-1} \|H\|_{\frac{r}{1-\varepsilon}} \\ &\leq \beta C \varepsilon \|\nabla u\|_r^{p-1} \|\nabla(\eta u)\|_r^{1-\varepsilon}. \end{aligned} \tag{3.9}$$

Notice that u plus a constant vector does not affect ∇u and the A -harmonic equation in (1.1) in our case, so let's assume that the average integral of u on $Q_{\frac{3}{2}R}$ is zero, and then by using Lemma 2.1,

$$\begin{aligned} \|\nabla(\eta u)\|_r^{1-\varepsilon} &= \|u \nabla \eta + \eta \nabla u\|_r^{1-\varepsilon} \\ &\leq (\|u \nabla \eta\|_r + \|\eta \nabla u\|_r)^{1-\varepsilon} \\ &\leq \left(\frac{C}{R} \|u\|_r + \|\nabla u\|_r\right)^{1-\varepsilon} \\ &\leq \left(\frac{C}{R} CR \|\nabla u\|_r + \|\nabla u\|_r\right)^{1-\varepsilon} \\ &\leq C \|\nabla u\|_r^{1-\varepsilon}, \end{aligned} \tag{3.10}$$

then we have

$$|I_1| \leq C \beta \varepsilon \|\nabla u\|_r^{1-\varepsilon}. \tag{3.11}$$

The estimate of I_2 is given below. By the hypothesis (H1), (3.5) and the definition of η , we have

$$\begin{aligned} |I_2| &= \left| \int_{Q_{\frac{3}{2}R}} \langle A(x, \nabla u), E(\eta, u) \rangle dx \right| \\ &\leq \beta \int_{Q_{\frac{3}{2}R}} |\nabla u|^{p-1} |E(\eta, u)| dx \end{aligned}$$

$$\begin{aligned} &\leq \beta 2^\varepsilon \frac{1+\varepsilon}{1-\varepsilon} \int_{Q_{\frac{3}{2}R}} |\nabla u|^{p-1} |u \nabla \eta|^{1-\varepsilon} dx \\ &\leq \beta 2^\varepsilon \frac{1+\varepsilon}{1-\varepsilon} C^{1-\varepsilon} \int_{Q_{\frac{3}{2}R}} |\nabla u|^{p-1} \left| \frac{u}{R} \right|^{1-\varepsilon} dx \\ &\leq \beta C \int_{Q_{\frac{3}{2}R}} |\nabla u|^{p-1} \left| \frac{u}{R} \right|^{1-\varepsilon} dx, \end{aligned} \tag{3.12}$$

By Young's inequality, for any $\theta_1 > 0$,

$$|I_2| \leq \beta C \theta_1 \int_{Q_{\frac{3}{2}R}} |\nabla u|^r dx + \beta C \int_{Q_{\frac{3}{2}R}} \left| \frac{u}{R} \right|^r dx. \tag{3.13}$$

For the second integral formula at the right end of the upper formula, take t such that $\max\{1, \frac{nr}{n+r}\} \leq t < r$, then by Lemma 2.1,

$$\beta C \int_{Q_{\frac{3}{2}R}} \left| \frac{u}{R} \right|^r dx \leq \beta C R^n \left(\int_{Q_{\frac{3}{2}R}} |\nabla u|^t dx \right)^{\frac{r}{t}}, \tag{3.14}$$

it doesn't effect ∇u and A -harmonic equation when u plus a constant, so assuming the integral average of u is zero in $Q_{\frac{3}{2}R}$, then we have

$$|I_2| \leq \beta C \theta_1 \int_{Q_{\frac{3}{2}R}} |\nabla u|^r dx + \beta C R^n \left(\int_{Q_{\frac{3}{2}R}} |\nabla u|^t dx \right)^{\frac{r}{t}}. \tag{3.15}$$

Combining the inequalities (3.7), (3.8), (3.11), (3.15), we obtain

$$\begin{aligned} \alpha \int_{Q_R} |\nabla u|^r dx &\leq C \beta \varepsilon \int_{Q_{\frac{3}{2}R}} |\nabla u|^r dx \\ &\quad + C \beta \theta_1 \int_{Q_{\frac{3}{2}R}} |\nabla u|^r dx \\ &\quad + C \beta R^n \left(\int_{Q_{\frac{3}{2}R}} |\nabla u|^t dx \right)^{\frac{r}{t}}. \end{aligned} \tag{3.16}$$

Divide the two sides of the formula above by $|Q_R| = \omega_n R^n$ (here ω_n is the volume of unit cube in \mathbf{R}^n), then

$$\begin{aligned} \alpha \int_{Q_R} |\nabla u|^r dx &\leq C \beta \varepsilon \int_{Q_{\frac{3}{2}R}} |\nabla u|^r dx + C \beta \theta_1 \int_{Q_{\frac{3}{2}R}} |\nabla u|^r dx \\ &\quad + C \beta \left(\int_{Q_{\frac{3}{2}R}} |\nabla u|^t dx \right)^{\frac{r}{t}}. \end{aligned} \tag{3.17}$$

Since Ω is bounded, $\Omega \subset Q_{R_0}$, $R < R_0$, the formula above becomes

$$\int_{Q_R} |\nabla u|^r dx \leq C(\varepsilon + \theta_1) \int_{Q_{\frac{3}{2}R}} |\nabla u|^r dx + C \left(\int_{Q_{\frac{3}{2}R}} |\nabla u|^r dx \right)^{\frac{r}{t}} \quad (3.18)$$

Let ε, θ_1 be small enough such that $\theta = C(\varepsilon + \theta_1) < 1$. Then

$$\int_{Q_R} |\nabla u|^r dx \leq \theta \int_{Q_{\frac{3}{2}R}} |\nabla u|^r dx + C \left(\int_{Q_{\frac{3}{2}R}} |\nabla u|^r dx \right)^{\frac{r}{t}}, \quad (3.19)$$

where $C = C(n, p, r, \alpha, \beta, R_0, \Omega)$. Noticing that the case we considered is that r is close enough to p , then r can be removed from the parameter of C . For $1 < t < r$, then (3.19) is a weak reverse Hölder inequality about ∇u .

Choosing $g = |\nabla u|^t, G = 0$ in $Q_{\frac{3}{2}R}$ and $g = G = 0$ in $Q_{2R} \setminus Q_{\frac{3}{2}R}$ with $q = \frac{r}{t}$. Then we arrive at the following inequality in $Q_{2R} \subset \Omega$, that is

$$\int_{Q_R} g^{\frac{r}{t}} dx \leq \theta \int_{Q_{2R}} g^{\frac{r}{t}} dx + C \left(\int_{Q_{2R}} g dx \right)^{\frac{r}{t}} + C \int_{Q_{2R}} G^{\frac{r}{t}} dx. \quad (3.20)$$

In the case (2), let $w = -\eta^p(u - u_0) \in W_0^{1,r}(Q_{2R})$, where $\eta \in C_0^\infty(Q_{2R})$ is a cutoff function, $0 \leq \eta \leq 1$, $|\nabla \eta| \leq \frac{C}{R}$, and $\eta \equiv 1$ in Q_R .

Extending the function $u - u_0$ with zero to $\mathbb{R}^n \setminus \Omega$ continuously. Then by Lemma 2.2, there exist $\phi \in W_0^{1, \frac{r}{1-\varepsilon}}(Q_{2R})$ and $H(x) \in L^{1-\varepsilon}(Q_{2R})$, such that

$$\begin{aligned} |\nabla w|^{-\varepsilon} \nabla w &= \nabla \phi + H \\ &= -|\nabla[\eta^p(u - u_0)]|^{-\varepsilon} \nabla[\eta^p(u - u_0)], \end{aligned} \quad (3.21)$$

and

$$\|H\|_{\frac{r}{1-\varepsilon}} \leq C\varepsilon \|\nabla[\eta^p(u - u_0)]\|_r^{1-\varepsilon}, \quad (3.22)$$

$$\|\nabla \phi\|_{\frac{r}{1-\varepsilon}} \leq C \|\nabla[\eta^p(u - u_0)]\|_r^{1-\varepsilon}, \quad (3.23)$$

where C is a constant only depending on n, r and Ω . Note that for Hodge decomposition (3.21), (3.22) and (3.23), we have $u - u_0 = 0$, H and $\nabla \phi$ are equal to zero when $u - u_0 \in \mathbb{R}^n \setminus \Omega$. By the Minkowski inequality and the selection of η , we have

$$\begin{aligned} \|\nabla[\eta^p(u - u_0)]\|_r^{1-\varepsilon} &\leq C \left[\|(u - u_0) \nabla \eta\|_r \right. \\ &\quad \left. + \|\nabla(u - u_0)\|_r \right]^{1-\varepsilon}. \end{aligned} \quad (3.24)$$

Noting that the boundary $\partial\Omega$ is r -Poincaré thick. Since $u - u_0$ is continually zero to $\mathbb{R}^n \setminus \Omega$, then by (1.2),

$$\begin{aligned} \|(u - u_0) \nabla \eta\|_r &\leq CR^{-1} \left(\int_{Q_{2R^1} \Omega} |u - u_0|^r dx \right)^{\frac{1}{r}} \\ &= CR^{-1} \left(\int_{Q_{2R}} |u - u_0|^r dx \right)^{\frac{1}{r}} \\ &\leq CR^{-1} \left(\int_{Q_{2R}} |\nabla(u - u_0)|^{\frac{nr}{n+r}} dx \right)^{\frac{n+r}{nr}} \\ &= CR^{-1} \left(\int_{Q_{2R^1} \Omega} |\nabla(u - u_0)|^{\frac{nr}{n+r}} dx \right)^{\frac{n+r}{nr}}, \end{aligned} \quad (3.25)$$

here we used $u - u_0 = 0$ in $\mathbb{R}^n \setminus \Omega$. Substitute the above formula into (3.24), we get

$$\begin{aligned} \|\nabla[\eta^p(u - u_0)]\|_r^{1-\varepsilon} &\leq C \left[R^{-1} \left(\int_{Q_{2R^1} \Omega} |\nabla(u - u_0)|^{\frac{nr}{n+r}} dx \right)^{\frac{n+r}{nr}} \right. \\ &\quad \left. + C \left(\int_{Q_{2R^1} \Omega} |\nabla(u - u_0)|^r dx \right)^{\frac{1}{r}} \right]^{1-\varepsilon}. \end{aligned} \quad (3.26)$$

Then (3.22) and (3.23) are

$$\begin{aligned} \left(\int_{Q_{2R} \cap \Omega} |H|^{\frac{r}{1-\varepsilon}} dx \right)^{\frac{1-\varepsilon}{r}} &\leq C\varepsilon \left[R^{-1} \left(\int_{Q_{2R} \cap \Omega} |\nabla(u - u_0)|^{\frac{nr}{n+r}} dx \right)^{\frac{n+r}{nr}} \right. \\ &\quad \left. + C \left(\int_{Q_{2R} \cap \Omega} |\nabla(u - u_0)|^r dx \right)^{\frac{1}{r}} \right]^{1-\varepsilon}, \end{aligned} \quad (3.27)$$

$$\begin{aligned} \left(\int_{Q_{2R^1} \Omega} |\nabla \phi|^{\frac{r}{1-\varepsilon}} dx \right)^{\frac{1-\varepsilon}{r}} &\leq C \left[R^{-1} \left(\int_{Q_{2R^1} \Omega} |\nabla(u - u_0)|^{\frac{nr}{n+r}} dx \right)^{\frac{n+r}{nr}} \right. \\ &\quad \left. + C \left(\int_{Q_{2R^1} \Omega} |\nabla(u - u_0)|^r dx \right)^{\frac{1}{r}} \right]^{1-\varepsilon}, \end{aligned} \quad (3.28)$$

By the conditions (H1), (H2), Lemma 2.4, Hodge decomposition (3.21), and the Definition 1.1, we obtain

$$\begin{aligned} &\alpha \int_{\Omega} \eta^{p(1-\varepsilon)} |\nabla u|^r dx \\ &\leq \int_{\Omega} \left\langle A(x, \nabla u), |\eta^p \nabla u|^{-\varepsilon} \eta^p \nabla u \right\rangle dx \\ &= \int_{\Omega} \left\langle A(x, \nabla u), |\eta^p \nabla u|^{-\varepsilon} \eta^p \nabla u \right. \\ &\quad \left. - |\eta^p \nabla(u - u_0)|^{-\varepsilon} \eta^p \nabla(u - u_0) \right\rangle dx \end{aligned}$$

$$\begin{aligned}
 & + \int_{\Omega} \left\langle A(x, \nabla u), |\eta^p \nabla(u - u_0)|^{-\varepsilon} \eta^p \nabla(u - u_0) \right. \\
 & \left. - |\nabla[\eta^p(u - u_0)]|^{-\varepsilon} \nabla[\eta^p(u - u_0)] \right\rangle dx \\
 & + \int_{\Omega} \left\langle A(x, \nabla u), |\nabla[\eta^p(u - u_0)]|^{-\varepsilon} \nabla[\eta^p(u - u_0)] \right\rangle dx \\
 & \leq C \int_{\Omega} \eta^{p(1-\varepsilon)} |A(x, \nabla u)| |\nabla u_0|^{1-\varepsilon} dx \\
 & + C \int_{\Omega} |A(x, \nabla u)| |\nabla \eta^p|^{1-\varepsilon} |u - u_0|^{1-\varepsilon} dx \quad (3.29) \\
 & - \int_{\Omega} \langle A(x, \nabla u), \nabla \phi + H \rangle dx \\
 & \leq C \beta \left[\int_{\Omega} \eta^{p(1-\varepsilon)} |\nabla u|^{p-1} |\nabla u_0|^{1-\varepsilon} dx \right. \\
 & + \int_{\Omega} \eta^{(p-1)(1-\varepsilon)} |\nabla u|^{p-1} |\nabla \eta|^{1-\varepsilon} |u - u_0|^{1-\varepsilon} dx \\
 & \left. + \int_{\Omega} |\nabla u|^{p-1} |H| dx \right] \\
 & @ C \beta [I_3 + I_4 + I_5].
 \end{aligned}$$

The estimate of I_3 is given below. By Young's inequality, for any $\theta_2 > 0$,

$$\begin{aligned}
 I_3 & = \int_{\Omega} \eta^{p(1-\varepsilon)} |\nabla u|^{p-1} |\nabla u_0|^{1-\varepsilon} dx \\
 & \leq \theta_2 \int_{\Omega} \eta^{r(1-\varepsilon)} |\nabla u|^r dx + C \int_{\Omega} \eta^r |\nabla u_0|^r dx \quad (3.30) \\
 & \leq \theta_2 \int_{Q_{2R^1 \Omega}} |\nabla u|^r dx + C \int_{Q_{2R^1 \Omega}} |\nabla u_0|^r dx.
 \end{aligned}$$

The estimate of I_4 is given below. By Young's inequality, for any $\theta_3 > 0$,

$$\begin{aligned}
 I_4 & = \int_{\Omega} \eta^{(p-1)(1-\varepsilon)} |\nabla u|^{p-1} |\nabla \eta|^{1-\varepsilon} |u - u_0|^{1-\varepsilon} dx \\
 & \leq \theta_3 \int_{\Omega} \eta^{r(1-\varepsilon)} |\nabla u|^r dx + C \int_{\Omega} |\nabla \eta|^r |u - u_0|^r dx \quad (3.31) \\
 & \leq \theta_3 \int_{Q_{2R^1 \Omega}} |\nabla u|^r dx + C \int_{Q_{2R^1 \Omega}} |\nabla \eta|^r |u - u_0|^r dx
 \end{aligned}$$

For the second integral formula at the right end of the upper formula, noticing that $\partial\Omega$ is r -Poincaré thick. By (3.25) we get

$$\begin{aligned}
 & C \int_{Q_{2R^1 \Omega}} |\nabla \eta|^r |u - u_0|^r dx \\
 & \leq CR^{-r} \left(\int_{Q_{2R^1 \Omega}} |\nabla(u - u_0)|^{\frac{nr}{n+r}} dx \right)^{\frac{n+r}{n}}. \quad (3.32)
 \end{aligned}$$

Then by Minkowski inequality and the Hölder inequality, we have

$$\begin{aligned}
 & \left(\int_{Q_{2R^1 \Omega}} |\nabla(u - u_0)|^{\frac{nr}{n+r}} dx \right)^{\frac{n+r}{n}} \\
 & \leq \left[\left(\int_{Q_{2R^1 \Omega}} |\nabla u|^{\frac{nr}{n+r}} dx \right)^{\frac{n+r}{n}} + \left(\int_{Q_{2R^1 \Omega}} |\nabla u_0|^{\frac{nr}{n+r}} dx \right)^{\frac{n+r}{n}} \right]^r \\
 & \leq \left[\left(\int_{Q_{2R^1 \Omega}} |\nabla u|^{\frac{nr}{n+r}} dx \right)^{\frac{n+r}{n}} + CR \left(\int_{Q_{2R^1 \Omega}} |\nabla u_0|^r dx \right)^{\frac{1}{r}} \right]^r
 \end{aligned}$$

$$\leq 2^r \left[\left(\int_{Q_{2R^1 \Omega}} |\nabla u|^{\frac{nr}{n+r}} dx \right)^{\frac{n+r}{n}} + CR^r \int_{Q_{2R^1 \Omega}} |\nabla u_0|^r dx \right], \quad (3.33)$$

Then

$$\begin{aligned}
 I_4 & \leq \theta_3 \int_{Q_{2R^1 \Omega}} |\nabla u|^r dx + CR^{-r} \left(\int_{Q_{2R^1 \Omega}} |\nabla u|^{\frac{nr}{n+r}} dx \right)^{\frac{n+r}{n}} \\
 & + C \int_{Q_{2R^1 \Omega}} |\nabla u_0|^r dx. \quad (3.34)
 \end{aligned}$$

The estimate of I_5 is given below. By Young's inequality, (3.27) and the Hölder inequality, for any $\theta_4 > 0$, we have

$$\begin{aligned}
 I_5 & = \int_{\Omega} |\nabla u|^{p-1} |H| dx \\
 & \leq \theta_4 \int_{\Omega} |\nabla u|^r dx + C \int_{\Omega} |H|^{\frac{r}{r-\varepsilon}} dx \\
 & \leq \theta_4 \int_{Q_{2R^1 \Omega}} |\nabla u|^r dx \\
 & + C \varepsilon \left[R^{-1} \left(\int_{Q_{2R^1 \Omega}} |\nabla(u - u_0)|^{\frac{nr}{n+r}} dx \right)^{\frac{n+r}{nr}} \right. \\
 & \left. + C \left(\int_{Q_{2R^1 \Omega}} |\nabla(u - u_0)|^r dx \right)^{\frac{1}{r}} \right]^r \quad (3.35)
 \end{aligned}$$

$$\begin{aligned}
 & \leq \theta_4 \int_{Q_{2R^1 \Omega}} |\nabla u|^r dx \\
 & + C \varepsilon \int_{Q_{2R^1 \Omega}} |\nabla(u - u_0)|^r dx \\
 & \leq \theta_4 \int_{Q_{2R^1 \Omega}} |\nabla u|^r dx + C \varepsilon \int_{Q_{2R^1 \Omega}} |\nabla u|^r dx \\
 & + C \varepsilon \int_{Q_{2R^1 \Omega}} |\nabla u_0|^r dx
 \end{aligned}$$

Combining the inequalities (3.29), (3.30), (3.34), (3.35), we obtain

$$\begin{aligned}
 & \int_{\Omega} \eta^{p(1-\varepsilon)} |\nabla u|^r dx \\
 & \leq C(\theta_2 + \theta_3 + \theta_4 + \varepsilon) \int_{Q_{2R^1 \Omega}} |\nabla u|^r dx \\
 & + CR^{-r} \left(\int_{Q_{2R^1 \Omega}} |\nabla u|^{\frac{nr}{n+r}} dx \right)^{\frac{n+r}{n}} \\
 & + C \int_{Q_{2R^1 \Omega}} |\nabla u_0|^r dx, \quad (3.36)
 \end{aligned}$$

where $C = C(n, p, \alpha, \beta, K, \Omega)$.

Choosing $\theta_2, \theta_3, \theta_4$ and $\varepsilon_0 > 0$ small enough, there exist $r_1 = p - \varepsilon_0 < p$, such that $\theta = C(\theta_2 + \theta_3 + \theta_4 + \varepsilon) < 1$ when $\varepsilon < \varepsilon_0$. By (3.36),

$$\begin{aligned}
 \int_{Q_R} |\nabla u|^r dx & \leq \theta \int_{Q_{2R}} |\nabla u|^r dx + C \left(\int_{Q_{2R}} |\nabla u|^r dx \right)^{\frac{r}{r_1}} \\
 & + C \int_{Q_{2R}} |\nabla u_0|^r dx, \quad (3.37)
 \end{aligned}$$

where $t = \frac{nr}{n+r} < r$. Let $g = |\nabla u|^r$, $G=0$. Then we arrive at the following inequality when $\varepsilon < \varepsilon_0$, that is

$$\int_{Q_R} g^{\frac{r}{t}} dx \leq \theta \int_{Q_{2R}} g^{\frac{r}{t}} dx + C \left(\int_{Q_{2R}} g dx \right)^{\frac{r}{t}} + C \int_{Q_{2R}} |G|^{\frac{r}{t}} dx, \quad (3.38)$$

where $C = C(n, p, \alpha, \beta, K, \Omega)$. Then by (3.20), (3.38) and Lemma 2.5, there exists r' , and $r' > r$, such that $u \in W^{1,r'}(\Omega)$. For r' , repeating the above process, the integrability of ∇u is improved over and over again. In this way, there must be an integrable exponent r_1 and r_2 , satisfying $r_1 < p < r_2$, such that $u \in W^{1,r}(\Omega)$, $\forall r \in (r_1, r_2)$. The proof is complete.

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